

Representation type of Jordan algebras

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Jordan algebras arose from the search of "exceptional" setting for quantum mechanics.

The "Copenhagen model":

- the physical observables are Hermitian matrices (or operators on Hilbert space);
- the basic operations are multiplication by $\lambda \in \mathbb{C}$, addition, multiplication of matrices, adjoint operator.

Most of matrix operations are not observable!

In 1932 the German physicist Pascual Jordan proposed a program to discover "*a new algebraic setting for quantum mechanics*", Jordan has chosen a new observable operation called **quasi-multiplication**

$$x \cdot y := \frac{1}{2}(xy + yx).$$

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Definition

A **Jordan algebra** J consists of a real vector space equipped with a bilinear product $x \cdot y$ satisfying the commutative law and the Jordan identity

$$\begin{aligned}x \cdot y &= y \cdot x, \\(x^2 \cdot y) \cdot x &= x^2 \cdot (y \cdot x), \quad \forall x, y \in J.\end{aligned}\tag{1}$$

A Jordan algebra is **formally real** if

$$x_1^2 + \cdots + x_n^2 = 0 \implies x_1 = \cdots = x_n = 0.$$

Any associative algebra A over \mathbb{R} gives rise to a Jordan algebra A^+ under quasi-multiplication: the product $x \cdot y$ is clearly commutative and satisfies the Jordan identity.

A Jordan algebra is called **special** if it can be realized as a Jordan subalgebra of some A^+ . For example, if A has an involution \star then the subspace of hermitian elements $H(A, \star) = \{x^\star = x \mid x \in A\}$ forms a special Jordan algebra.

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These hermitian algebras are the archetypes of all Jordan algebras. It is easy to check that the hermitian matrices over \mathbb{R} , \mathbb{C} and the quaternions form special Jordan algebras that are formally real.

One obtains another special formally real Jordan algebra (called **spin factor** $JSpin_n$ or **Jordan algebra of Clifford type**) on the space $\mathbb{R}1 \oplus \mathbb{R}^n$, when $n \geq 2$, by making 1 acts as a unit and defining $v \cdot w = \langle v, w \rangle 1$, where $v, w \in \mathbb{R}^n$ and $\langle \cdot, \cdot \rangle$ is an inner product in \mathbb{R}^n .

What the physicists were looking for were Jordan algebras where there is no invisible structure xy governing the visible structure $x \cdot y$. A Jordan algebra \mathcal{J} is called **exceptional** if it is not special, i.e. does not result from quasi-multiplication.

Jordan hoped that by studying finite-dimensional Jordan algebras he could find families of simple exceptional algebras E_n , parametrized by natural numbers n , letting n to infinity would provide a suitable infinite-dimensional home for quantum mechanics.

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In 1934 Jordan, John von Neumann and Eugene Wigner showed that every finite-dimensional formally real Jordan algebra can be written as a direct sum of so-called simple ones, which are four infinite families, together with one exceptional case:

- The Jordan algebra of $n \times n$ self-adjoint real matrices, complex matrices and quaternionic matrices;
- The Jordan algebra of Clifford type corresponding to usual inner product in \mathbb{R}^n ;
- The Jordan algebra $H_3(\mathbb{O})$ of 3×3 self-adjoint octonionic matrices, (an exceptional Jordan algebra called the **Albert algebra**, since A.Albert showed that it is indeed exceptional).

There was only one exceptional algebra in the list!!! $H_n(\mathbb{O})$, $n > 3$ is not Jordan, while $H_2(\mathbb{O}) \simeq JSpin_9$. The lone $H_3(\mathbb{O})$ was too isolated to give a clue as to possible existence of inf-dimensional exceptional algebras.

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Let $\mathbf{k} = \bar{\mathbf{k}}$, $\text{char } \mathbf{k} = 0$. A **Jordan algebra** is a commutative \mathbf{k} -algebra (J, \cdot) satisfying Jordan identity

$$(x^2 \cdot y) \cdot x = x^2 \cdot (y \cdot x) \quad x, y \in J. \quad (2)$$

Example

- 1 Denote $H_n(A) = H(M_n(A), \bar{j}) \subset M_n(A)^+$.
- 2 Let f be a symmetric bilinear form on E , then $J(E, f) = E \oplus \mathbf{k}$ with respect to

$$(\alpha + x) \cdot (\beta + y) = \alpha\beta + f(x, y)1 + (\alpha y + \beta x),$$

$\alpha, \beta \in \mathbf{k}$, $x, y \in E$ is a Jordan algebra of Clifford type.

- 3 Let \mathbb{O} be an Cayley algebra over \mathbf{k} , then $\mathcal{A} = H_3(\mathbb{O})$ is called the **Albert algebra**

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Suppose M is a \mathbf{k} -vector space equipped with two bilinear mappings

$$l, r : J \otimes M \rightarrow M \quad l : (a, m) \rightarrow am, \quad r : (a, m) \rightarrow ma,$$

and define a product on $\Omega = J \oplus M$

$$(a_1 + m_1) \circ (a_2 + m_2) = a_1 \cdot a_2 + a_1 m_2 + m_1 a_2.$$

M is a **Jordan bimodule** for $J \Leftrightarrow \Omega = (\Omega, \circ)$ is a Jordan algebra.

Equivalently a linear map $\rho : J \rightarrow \text{End}_k M$, $\rho(a)m = am$ defines a **(bi)representation** if for all $a, b \in J$

$$[\rho(a), \rho(a \cdot a)] = 0,$$

$$2\rho(a)\rho(b)\rho(a) + \rho((a \cdot a) \cdot b) = 2\rho(a)\rho(a \cdot b) + \rho(b)\rho(a \cdot a).$$

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For any J the category of \mathbf{k} -finite dimensional J -bimodules is equivalent to the category of finite dimensional left U -modules for some associative algebra called the **universal multiplicative enveloping algebra** $U = U(J)$.

The algebra U can be constructed as a quotient of the tensor algebra $F(J)/R$, where R is an ideal generated by the elements:

$$\begin{aligned} & a_1 \otimes a_2 \otimes a_3 + a_3 \otimes a_2 \otimes a_1 + (a_1 \cdot a_3) \cdot a_2 \\ & \quad - a_1 \otimes a_2 \cdot a_3 - a_2 \otimes a_1 \cdot a_3 - a_3 \otimes a_1 \cdot a_2, \\ & a_1 \otimes a_2 \cdot a_3 + a_2 \otimes a_1 \cdot a_3 + a_3 \otimes a_1 \cdot a_2 \\ & - a_2 \cdot a_3 \otimes a_1 - a_1 \cdot a_3 \otimes a_2 - a_1 \cdot a_2 \otimes a_3, \quad a_1, a_2, a_3 \in J. \end{aligned}$$

We denote by $J\text{-mod}$ the category of finite dimensional J -bimodules.

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Jordan algebras versus Lie algebras

- 1 Let A be an associative algebra we construct $A^+ = (A, \frac{1}{2}(ab + ba))$, $A^- = (A, ab - ba)$.
- 2 For $\forall \mathfrak{g} \dim_{\mathbb{k}} U(\mathfrak{g}) = \infty$ while $\dim_{\mathbb{k}} J < \infty \Rightarrow \dim_{\mathbb{k}} U(J) < \infty$.
- 3 For any Lie algebra \mathfrak{g} its universal enveloping algebra $U(\mathfrak{g}) = F(\mathfrak{g})/\langle ab - ba - [a, b] \rangle$ and by PBW thm

$$\mathfrak{g} \subset U(\mathfrak{g})^- \text{ e } \mathfrak{g}\text{-mod} \simeq U(\mathfrak{g})\text{-mod.}$$

- 4 $S(J) = F(J)/\langle ab + ba - 2a \cdot b \rangle$ is called **universal associative enveloping** of J and

$$J \subset S(J)^+ \iff J \text{ is special, } J\text{-mod}_{\frac{1}{2}} \simeq S(J)\text{-mod.}$$

Let $\sigma : J \rightarrow \text{End}_{\mathbb{k}} M$, satisfying for all $a, b \in J$

$$\sigma(a \cdot b) = \frac{1}{2}(\sigma(a)\sigma(b) + \sigma(b)\sigma(a)). \quad (3)$$

This gives the structure of J -bimodule on M via $\frac{1}{2}\sigma : J \rightarrow \text{End}_{\mathbb{k}} M$, and is called **one-sided (bi)modules**.

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Let $\rho : a \rightarrow \rho_a$ is a representation of J corresponding to M then

$$\rho_e(\rho_e - 1)(2\rho_e - 1) = 0$$

gives the Peirce decomposition for M

$$M = M_0 \oplus M_{\frac{1}{2}} \oplus M_1,$$

where $M_i = \{m_i \mid \rho_e m_i = i m_i\}$.

In particular $\rho : J \rightarrow \text{End}_{\mathbf{k}} M_1$, is a unital representation, and $\rho : J \rightarrow \text{End}_{\mathbf{k}} M_{\frac{1}{2}}$ is one-sided representation of J , while $\rho : J \rightarrow \text{End}_{\mathbf{k}} M_0$ is one-dimensional trivial representation.

Therefore to describe $J\text{-mod}$ it suffices to describe $S(J)\text{-mod}$ and $U_1(J)\text{-mod}$.

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We start with the Albert classification of simple finite-dimensional:

- \mathbf{k} ;
- $J(V, f)$, f non-degenerate;
- $H_n(C)$, $n \geq 3$, (C, τ) composition algebra of dimension 1, 2, 4;
- \mathcal{A} , the Albert algebra.

In 1954, N. Jacobson described all irreducible modules for J simple:

J	$S(J)$	$U_1(J)$
$Sym_n(\mathbf{k})$	M_n	$M_{\frac{n(n+1)}{2}} \oplus M_{\frac{n(n-1)}{2}}$
$M_n(\mathbf{k})^+$	$M_n \oplus M_n$	$M_{n^2} \oplus M_{\frac{n(n+1)}{2}} \oplus M_{\frac{n(n+1)}{2}}$ $\oplus M_{\frac{n(n-1)}{2}} \oplus M_{\frac{n(n-1)}{2}}$
$Symp_{2n}(\mathbf{k})$	M_{2n}	$M_{n(2n-1)} \oplus M_{n(2n+1)}^*$
$J_n(V, f)$ $\dim V = n$ is even	M_{2^n}	$\oplus_s M_s$ $s = \binom{n+1}{1}, \binom{n+1}{3}, \dots, \binom{n+1}{n+1}$
$J_n(V, f)$ $n = 2\nu - 1$	$M_{2^{n-1}} + M_{2^{n-1}}$	$M_{\frac{1}{2}\binom{n+1}{\nu}} \oplus M_{\frac{1}{2}\binom{n+1}{\nu}} \oplus_s M_s$ $s = \binom{n+1}{0}, \binom{n+1}{1}, \dots, \binom{n+1}{\nu-1}$
\mathcal{A}	0	$M_{27}(\mathbf{k})$

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$Symp_{2n}(\mathbf{k})$	M_{2n}	$M_{n(2n-1)} \oplus M_{n(2n+1)}^*$
$J_n(V, f)$ $\dim V = n$ is even	M_{2^n}	$\oplus_s M_s$ $s = \binom{n+1}{1}, \binom{n+1}{3}, \dots, \binom{n+1}{n+1}$
$J_n(V, f)$ $n = 2\nu - 1$	$M_{2^{n-1}} + M_{2^{n-1}}$	$M_{\frac{1}{2}\binom{n+1}{\nu}} \oplus M_{\frac{1}{2}\binom{n+1}{\nu}} \oplus_s M_s$ $s = \binom{n+1}{0}, \binom{n+1}{1}, \dots, \binom{n+1}{\nu-1}$
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We start with the Albert classification of simple finite-dimensional:

- \mathbf{k} ;
- $J(V, f)$, f non-degenerate;
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In 1954, N. Jacobson described all irreducible modules for J simple:

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Let S_1 and S_2 be in $J\text{-mod}_{\frac{1}{2}}$ then we may form the **Kronecker sum** of two given one-sided modules $S = S_1 \otimes S_2 \in J\text{-mod}_1$ by setting

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Representation type of algebra

Let k be algebraically closed field and A be an associative finite dimensional k -algebra. Then A is of

- a **finite representation type** if there are finitely many isomorphism classes of finitely generated, indecomposable left A -modules.
- a **tame representation type** if isomorphism classes of finitely generated, indecomposable left A -modules form in each dimension finitely many one-parameter families.
- or a **wild representation type**

$A \iff$ quiver with relations $(Q(A), R)$ such that

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Question: What class of algebras to choose?

Inspired by results of P. Gabriel for associative case:

$$J = J_s + \text{Rad}J, \quad \text{Rad}^2 J = 0$$

Let Q be a quiver, the **quiver double** $D(Q)$ of Q is defined as follows:

$$D(Q_0) = \{X^+, X^- \mid X \in Q_0\}$$

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Let A be a finite dimensional associative algebra over algebraically closed field, such that $\text{Rad}^2 A = 0$, Q its quiver. Then A is of finite (tame) representation type if and only if $D(Q)$ is a disjoint union of simply-laced Dynkin diagrams (extended Dynkin diagrams).

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- 1 describe **matrix** Jordan algebras which are tame/finite with respect to their **one-sided** representation type:
K., Ovsienko S., Shestakov I., Representation type of Jordan algebras, Advances in Math., 2011.
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The Tits-Kantor-Koecher construction

A **short grading** of \mathfrak{g} is a \mathbb{Z} -grading of the form $\mathfrak{g} = \mathfrak{g}_{-1} \oplus \mathfrak{g}_0 \oplus \mathfrak{g}_1$.

Let P be the commutative bilinear map on J : $P(x, y) = x \cdot y$.

We associate to J a Lie algebra with short grading

$$\text{Lie}(J) = \mathfrak{g}_{-1} \oplus \mathfrak{g}_0 \oplus \mathfrak{g}_1.$$

Put $\mathfrak{g}_{-1} = J$, $\mathfrak{g}_0 = \langle L_a, [L_a, L_b] \mid a, b \in J \rangle$, $\mathfrak{g}_1 = \langle P, [L_a, P] \mid a \in J \rangle$

- $[L, x] = L(x)$ for $x \in \mathfrak{g}_{-1}$, $L \in \mathfrak{g}_0$;
- $[B, x](y) = B(x, y)$ for $B \in \mathfrak{g}_1$ and $x, y \in \mathfrak{g}_{-1}$;
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Then $\mathfrak{g} = \text{Lie}(J)$ is Lie algebra and is called the **Tits-Kantor-Koecher (TKK) construction** for J .

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A **short subalgebra** of \mathfrak{g} is an \mathfrak{sl}_2 subalgebra spanned by e, h, f such that the eigenspace decomposition of $ad h$ defines a short grading on \mathfrak{g} .

For any J with identity e consider in $Lie(J)$

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Let $\mathfrak{g} = \mathfrak{g}_{-1} \oplus \mathfrak{g}_0 \oplus \mathfrak{g}_1$ be the \mathbb{Z}_2 -graded Lie algebra, $p \in \mathfrak{g}_1$. For any $x, y \in \mathfrak{g}_{-1}$ set

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then $Jor(\mathfrak{g}) := (\mathfrak{g}_{-1}, \cdot)$ is a Jordan algebra.

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Relations between J -mod and $\mathfrak{g} = TKK(J)$ -modules?

We define two adjoint functors Jor and Lie between J -mod and \mathfrak{g} -modules admitting a short grading.

Not every J -module can be obtained from a \mathfrak{g} -module by application of Jor : one has to consider $\hat{\mathfrak{g}}$ the universal central extension of \mathfrak{g} .

Let \mathcal{S} (resp. $\mathcal{S}_{\frac{1}{2}}$) be the category of $\hat{\mathfrak{g}}$ -modules M such that the action of α_J induces a short grading on M (resp. a grading of length 2, namely $M_{-\frac{1}{2}} \oplus M_{\frac{1}{2}}$).

To define Jor let $N \in \mathcal{S}$. Then $N = N_1 \oplus N_0 \oplus N_{-1}$. We set $Jor(N) := N_{-1}$

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It is clear that Jor is an exact functor.

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where N is an abelian Lie algebra and $N_{-1} = M$.
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Let $\mathcal{P} = \hat{\mathfrak{g}}_0 \oplus \mathfrak{g}_{-1}$ and we extend the above $\hat{\mathfrak{g}}_0$ -module structure on M to a \mathcal{P} -module structure by setting $\mathfrak{g}_{-1}M = 0$. Next

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Let $\mathcal{P} = \hat{\mathfrak{g}}_0 \oplus \mathfrak{g}_{-1}$ and we extend the above $\hat{\mathfrak{g}}_0$ -module structure on M to a \mathcal{P} -module structure by setting $\mathfrak{g}_{-1}M = 0$. Next

$$\Gamma(M) = U(\hat{\mathfrak{g}}) \otimes_{U(\mathcal{P})} M.$$

We define $Lie(M)$ to be the maximal quotient in $\Gamma(M)$ which belongs to \mathcal{S} .

- $Jor \circ Lie$ is isomorphic to the identity functor in $J - mod_1$.
- Let $N \in \mathcal{S}$ and $\hat{\mathfrak{g}}N = N$, then the canonical map $Lie(Jor(N)) \rightarrow N$ is surjective.
- Let $N \in \mathcal{S}$ and $N^{\hat{\mathfrak{g}}} := \{x \in N \mid \hat{\mathfrak{g}}x = x\} = 0$, then the canonical map $N \rightarrow Lie(Jor(N))$ is injective.
- If $M \rightarrow L \rightarrow 0$ is exact in $J - mod_1$, then $Lie(M) \rightarrow Lie(L) \rightarrow 0$ is exact in \mathcal{S} .

The splitting $J - mod_1 \oplus J - mod_0$ can not be lifted to the Lie algebra $\hat{\mathfrak{g}}$, since some modules can have non-trivial extensions with trivial modules, thus left and right adjoint of the functor Jor are not isomorphic and the categories \mathcal{S} and $J - mod_1$ are not equivalent. Still they are close enough and one can describe projective modules, quivers and relations of $J - mod_1$ in terms of \mathcal{S} .

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Representation table for $TKK(J)$

J	\mathfrak{g}	$\mathcal{S}_{\frac{1}{2}}$	\mathcal{S}
$Sym_n(\mathbf{k})$	\mathfrak{sp}_{2n}	V	$ad, \Lambda^2 V$
$M_n(\mathbf{k})^+$	\mathfrak{sl}_{2n}	V, V^*	$ad, S^2(V), S^2(V^*), \Lambda^2(V), \Lambda^2(V^*)$
$Symp_{2n}(\mathbf{k})$	\mathfrak{so}_{4n}	V	$ad, S^2(V)$
$J_n(E, f)$ $n = 2\nu$	\mathfrak{so}_{n+3}	Γ spinor	$\Lambda^i(V), i = 1, \dots, \nu + 1$
$J_n(E, f)$ $n = 2\nu - 1$	\mathfrak{so}_{n+3}	Γ^+, Γ^- spinor	$\Lambda^i(V), i = 1, \dots, \nu$ $\Lambda^{\nu+1}(V)^\pm$
\mathcal{A}	E_7		ad

Quiver of an abelian category

Let \mathcal{C} be an abelian category with finitely many simple modules such that every object has finite length and every simple object has a projective cover.

Then \mathcal{C} is equivalent to the category of finite-dimensional A -modules. If L_1, \dots, L_r is the set of all up to isomorphism simple objects in \mathcal{C} and P_1, \dots, P_r are their projective covers, then A is a pointed algebra which is usually realized as the path algebra of a certain quiver Q with relations.

The vertices

$$Q_0 = \{\text{simple modules } L_1, \dots, L_r\}$$

$$Q_1 = \{\#\text{arrows from vertex } L_j \text{ to vertex } L_i \text{ is } \dim \text{Ext}^1(L_j, L_i)\}$$

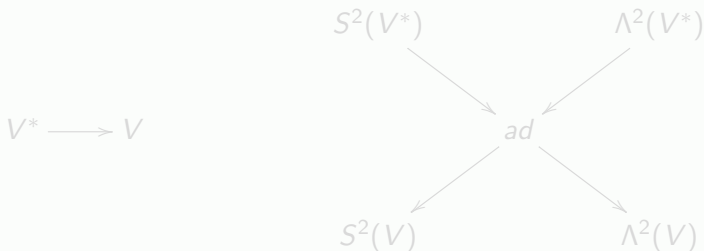
It is now clear how to describe the quiver Q .

Lemma

Let $\mathfrak{g} = \mathfrak{g}_s + \mathfrak{r}$ be the Levi decomposition of \mathfrak{g} . Let L and L' be two simple \mathfrak{g}_s -modules then $\dim \text{Ext}^1(L, L')$ equals the multiplicity of L' in $L \otimes \mathfrak{r}$.

Example

Let $J = M_n^{(+)}(\mathbb{k}) + \text{Sym}_n$ then its $\text{TKK}(J) = \mathfrak{sl}_{2n} + S^2(V)$, while $Q_{S^{\frac{1}{2}}}(J)$ and $Q_S(J)$ are correspondingly



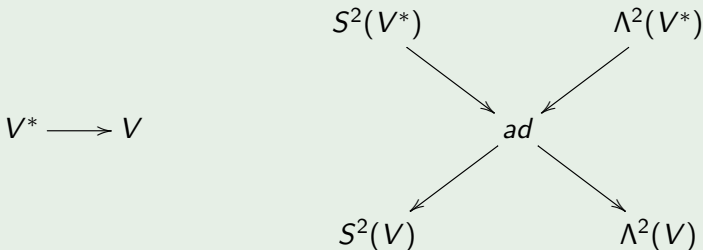
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Example

Let $J = M_n^{(+)}(\mathbf{k}) + \text{Sym}_n$ then its $\text{TKK}(J) = \mathfrak{sl}_{2n} + S^2(V)$, while $Q_{S^2}(\frac{1}{2}(J))$ and $Q_S(J)$ are correspondingly



Example

$$\mathfrak{g} = \mathfrak{so}_{2m+1} \oplus V, \quad m \geq 3$$

$$tr \begin{array}{c} \xrightarrow{\gamma_0} \\ \xleftarrow{\delta_0} \end{array} V \begin{array}{c} \xrightarrow{\gamma_1} \\ \xleftarrow{\delta_1} \end{array} \Lambda^2 V \begin{array}{c} \xrightarrow{\gamma_2} \\ \xleftarrow{\delta_2} \end{array} \dots \begin{array}{c} \xrightarrow{\gamma_{m-1}} \\ \xleftarrow{\delta_{m-1}} \end{array} \Lambda^{m-1} V \begin{array}{c} \xrightarrow{\gamma_m} \\ \xleftarrow{\delta_m} \end{array} \Lambda^m V \end{array}$$

with the relations

$$\begin{aligned} \gamma_{r-1}\gamma_r &= \delta_r\delta_{r-1} = 0, \\ \gamma_{r-1}\delta_{r-1} &= \delta_r\gamma_r, \\ \gamma_{m-1}\delta_{m-1} &= \gamma_m^2, \\ \text{for } r &= 1, \dots, m-1. \end{aligned}$$